

JF Semblat

Seismicity

Seismic waves

Elastodyn.

Amplification

Basin effects

Attenuation

Viscoelastic approxim.

Damped waves

Perspectives

GEOMECHANICS AND ADVANCED GEOTECHNICS

Engineering seismology and soil dynamics.

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Introduction

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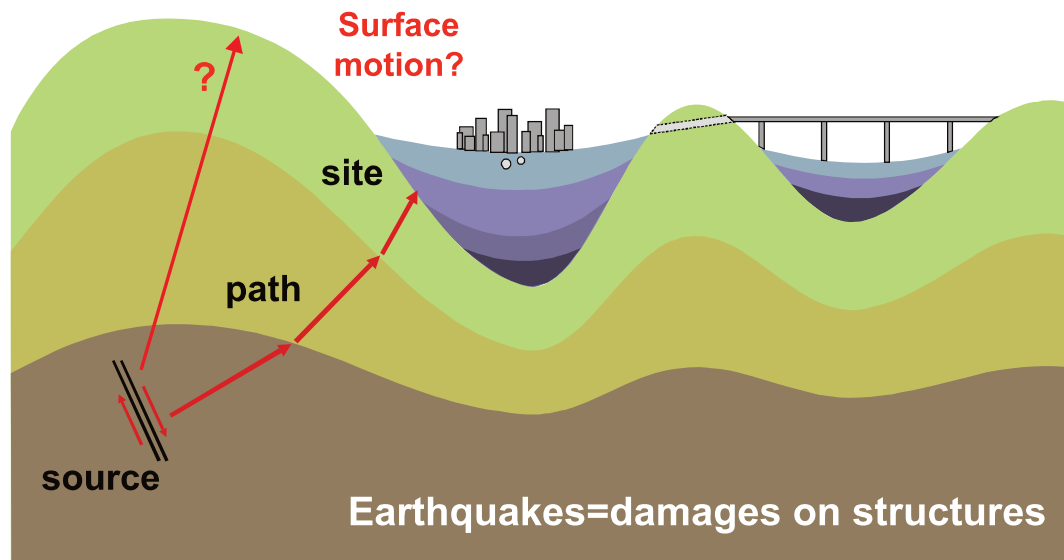
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- From source to structures through site:



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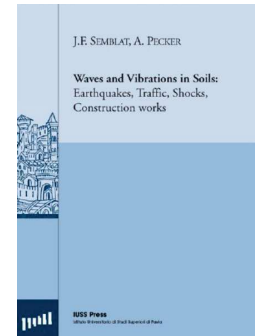
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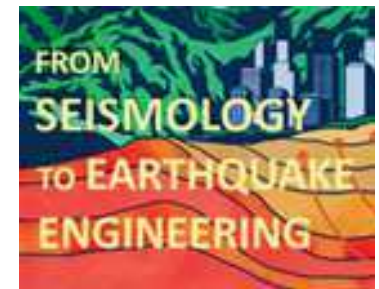
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■ References:

- *Waves and Vibrations in Soils: Earthquakes, Traffic, Shocks, Construction Works*, Semblat J-F, Pecker A., IUSS Press, 2009.



- MOOC "From seismology to earthquake engineering":
Teaser: youtu.be/V6mnVBszk0E
Enroll: www.coursera.org



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Map and fault (GNS)



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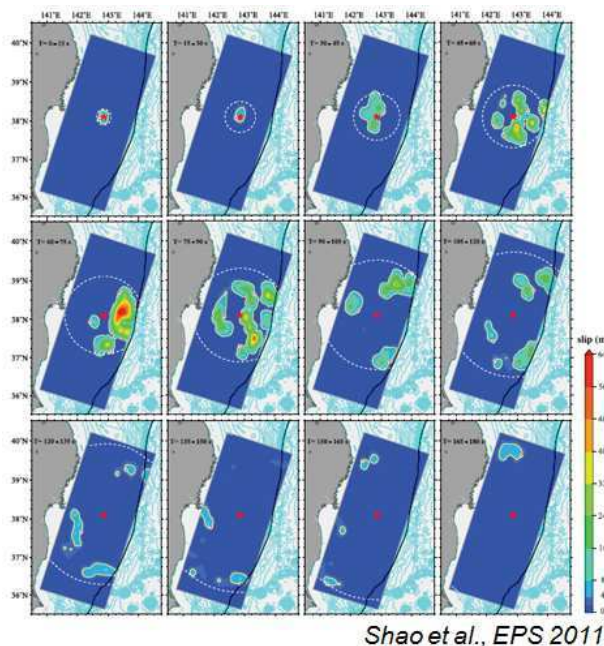
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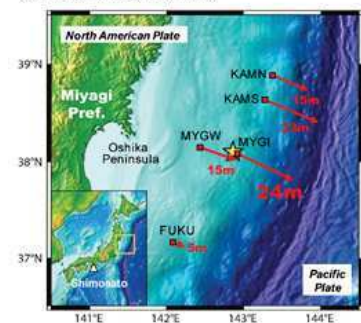
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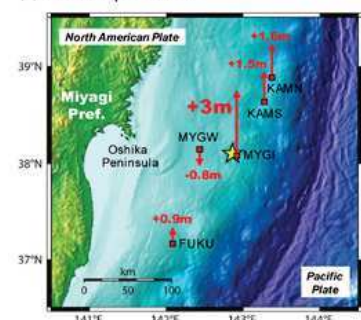
Displacements at the fault & at the surface



(A) Horizontal displacements



(B) Vertical displacements



Sato et al., Science 2011

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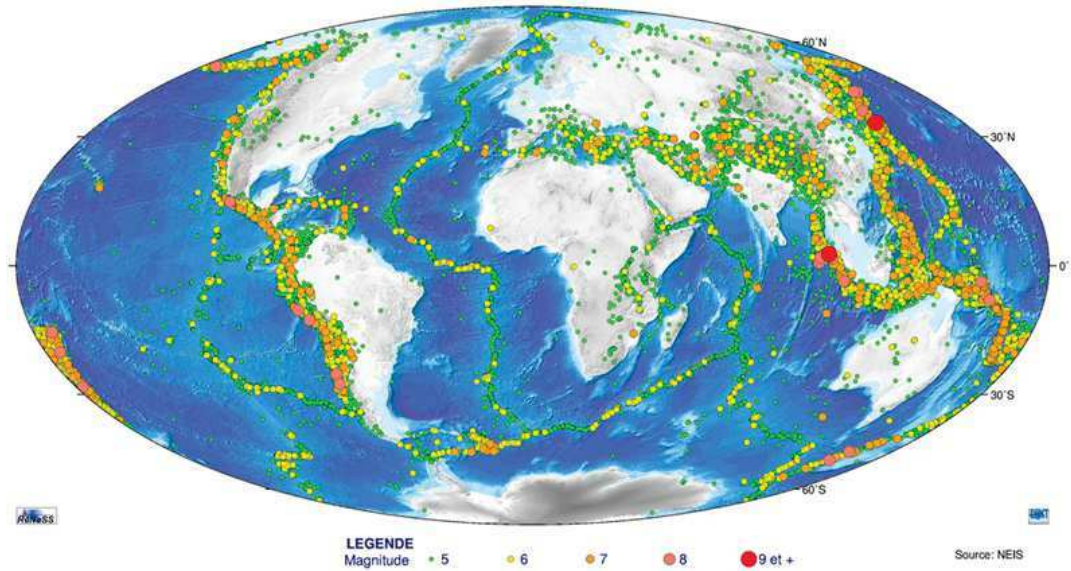
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■ Earthquakes in the world from 1973 to 2012



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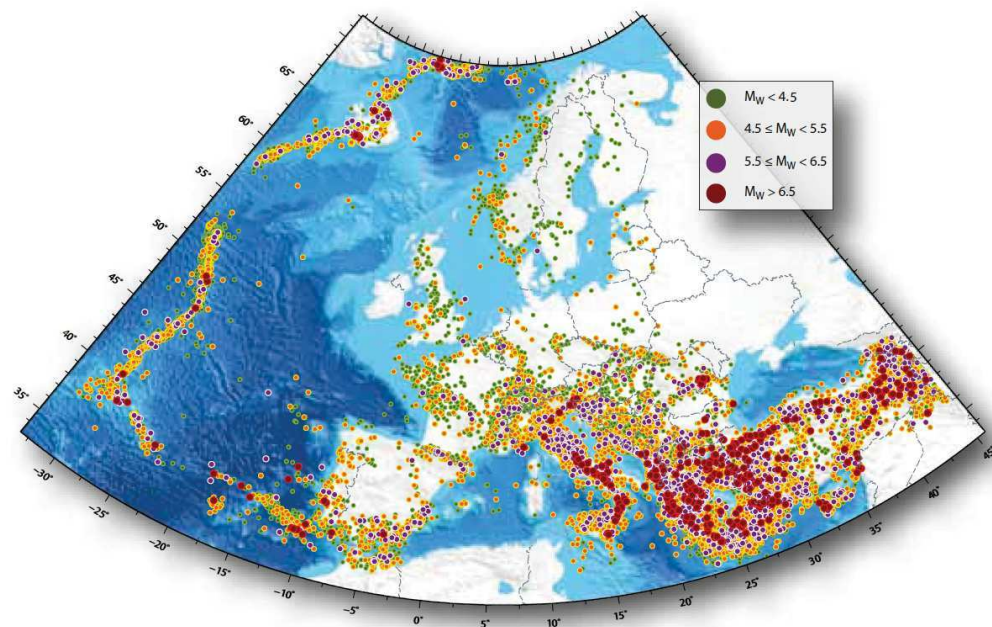
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■ Earthquakes in Europe for the period 1000-2007 with moment magnitudes $M_w \geq 3.5$ (SHARE).



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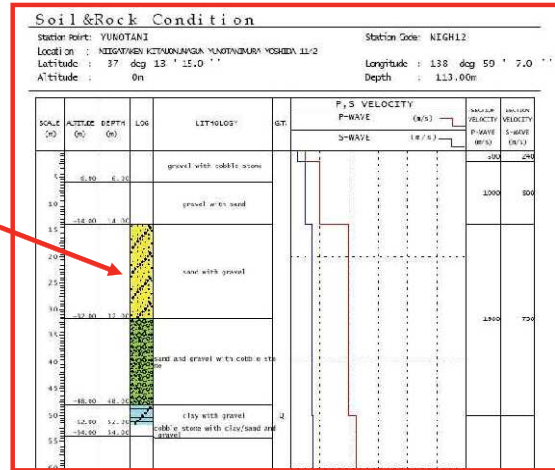
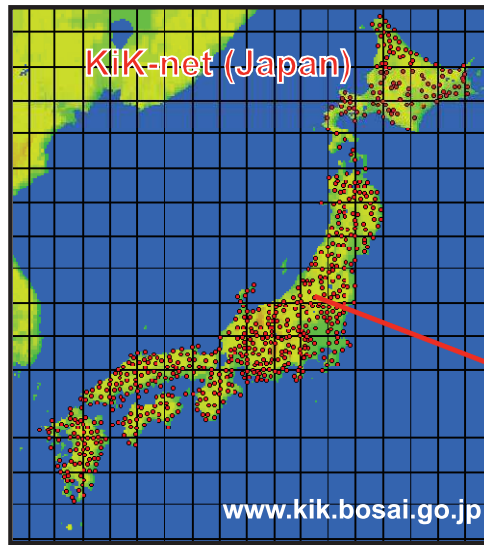
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■ Kik-net: observations in Japan



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- Uncoupled elastic wave equations:

$$\begin{cases} \Delta(\phi) = \frac{1}{V_P^2} \frac{\partial^2 \phi}{\partial t^2} \text{ with } V_P = \sqrt{\frac{\lambda + 2\mu}{\rho}} \\ \Delta(\psi) = \frac{1}{V_S^2} \frac{\partial^2 \psi}{\partial t^2} \text{ with } V_S = \sqrt{\frac{\mu}{\rho}} \end{cases}$$

- Typical velocity values:

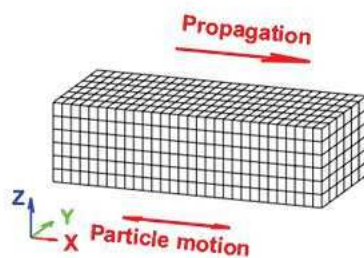
Material	V_S (m/s)	V_P (m/s)
Sand	200-400	400-800
Clay	100-200	1500
Rock	>800	>2000

$$\frac{V_P}{V_S} = \sqrt{\frac{2 - 2\nu}{1 - 2\nu}}$$

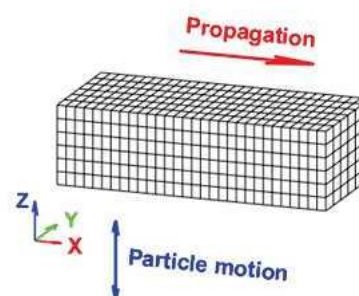
- Disregarding the boundary conditions, the wave equations define body waves as follows (Poisson, 1828):

- primary or pressure waves (P-waves)
- secondary or shear waves (S-waves)

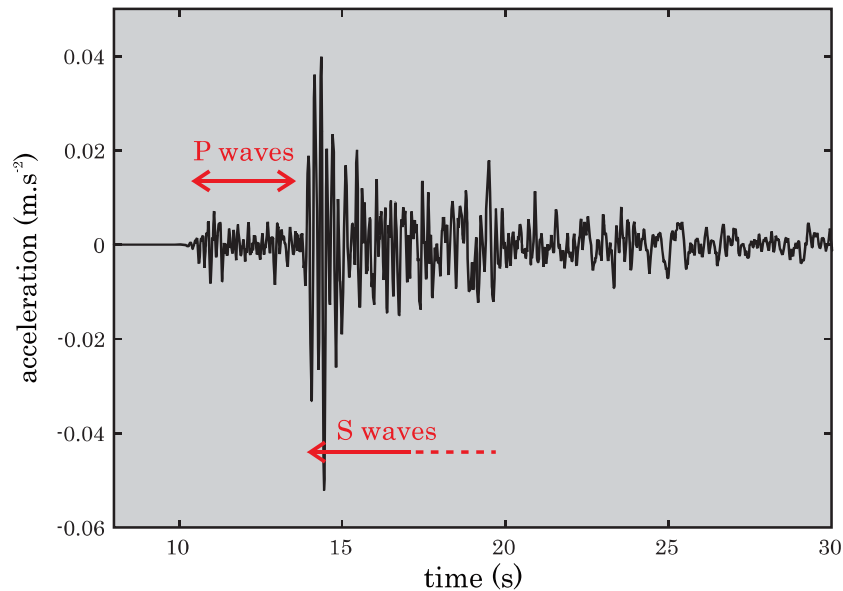
P wave



SV wave

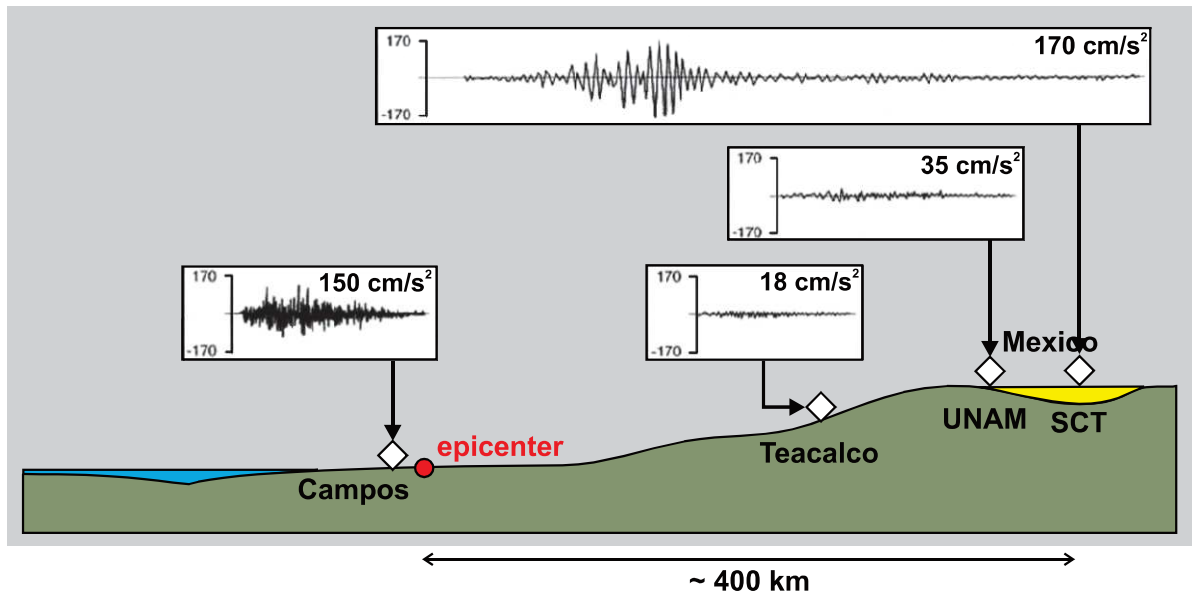


- In 1897, Richard D. Oldham identifies the *different wave types* in a seismic recording
- Ex.: earthquake recording in the city of Nice (CEREMA)



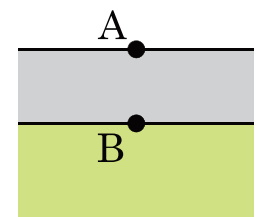
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■ Mexico-Michoacan earthquake (M8.0, sept., 19, 1985)



■ Elastic layer on an elastic half-space, harmonic SH-wave: spectral ratio A/B :

$$T(\omega) = \frac{u_A}{u_B} = \frac{2A_1}{A_2 + A'_2} \Rightarrow |T(\omega)| = \frac{1}{\cos(k_{z_1} h)}$$



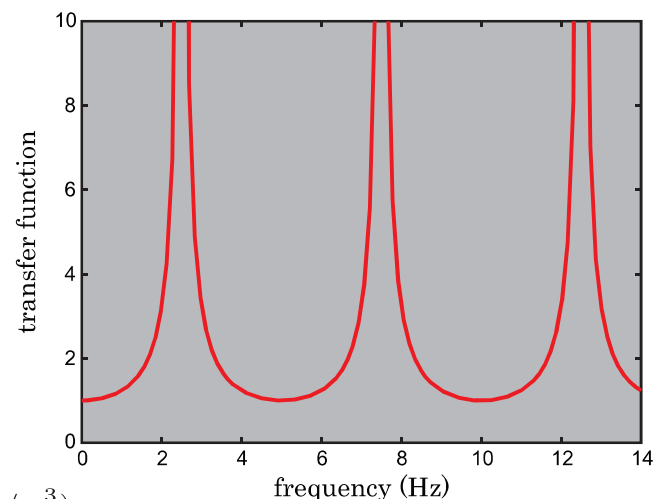
independent of μ_2 !

■ $T(\omega)$ infinite for:

$$\omega_n = \frac{(2n-1)\pi}{2} \frac{V_{S_1}}{h \cos \theta_1}$$

$$\omega_1 = 2\pi f_1 \Rightarrow$$

$$f_1 = \frac{V_{S_1}}{4h} \frac{1}{\cos \theta_1}$$



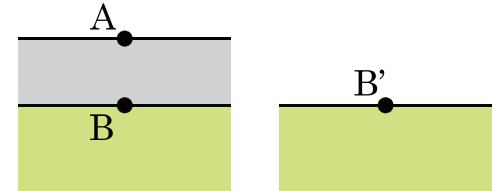
($h=20\text{m}$, $V_{S_1}=200\text{m/s}$, $\rho=2000\text{kg/m}^3$)

- Alternative spectral ratio: motions at A and B'
- B' : hypothetical location at interface without layer!

$$T^*(\omega) = \frac{u_A}{u_{B'}} = \frac{2A_1}{2A_2} \Rightarrow$$

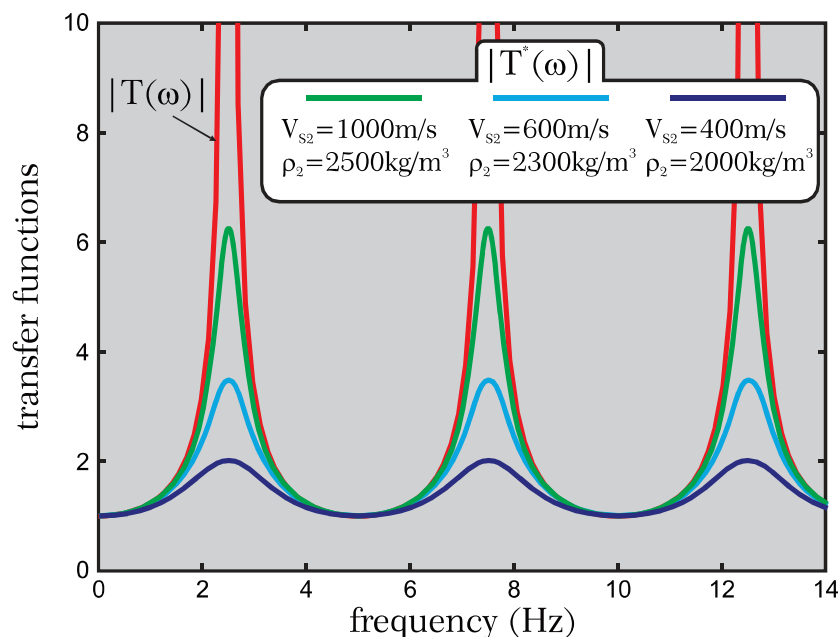
$$T^*(\omega) = \frac{1}{\cos(k_{z_1}h) + i\bar{\chi} \sin(k_{z_1}h)}$$

$$\text{with: } \bar{\chi} = \sqrt{\frac{\mu_1 \rho_1 \cos \theta_1}{\mu_2 \rho_2 \cos \theta_2}}$$



- $T^*(\omega) \neq T(\omega)$ since:
 - It depends on the properties of both media ($\bar{\chi}$)
 - The peaks have a finite amplitude $= 1/\bar{\chi}$: energy partially absorbed by the reflected wave A'_2 (radiative damping). It is not the case when $\bar{\chi}$ is large.

- Amplification for various velocity contrasts



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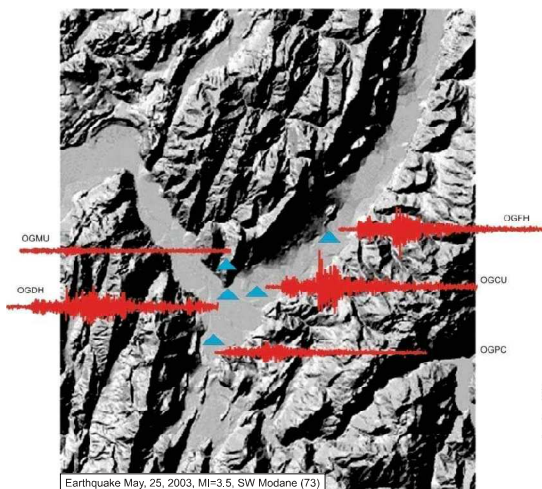
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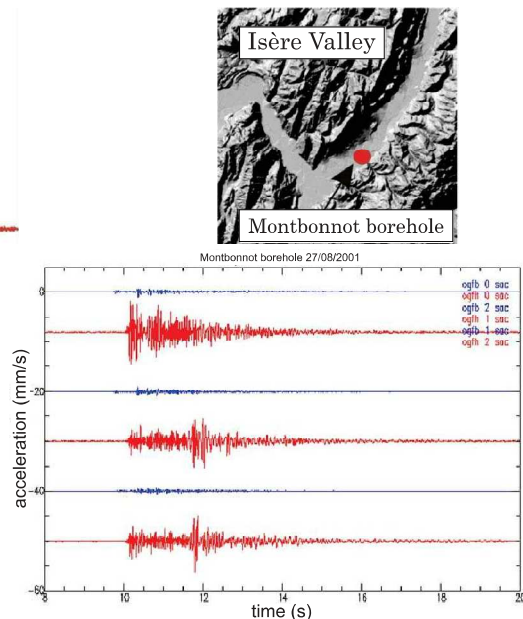
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■ Deep alpine valley (Grenoble, France)

Site effects in Grenoble

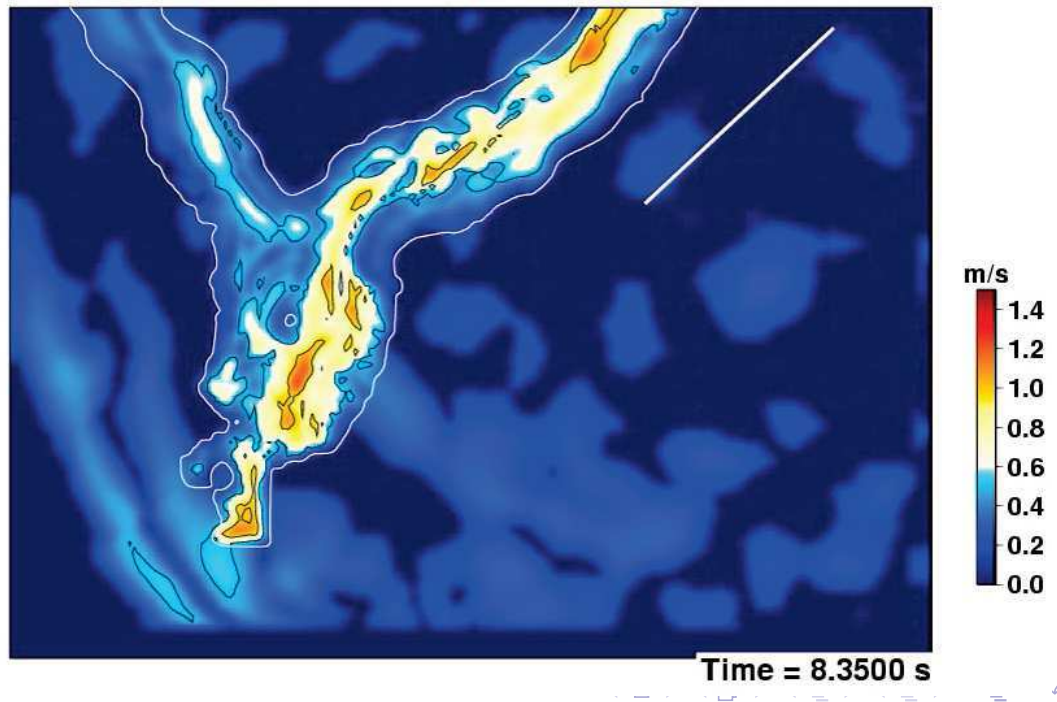


Montbonnot borehole



- Deep alpine valley (Grenoble, France): SEM simulations

Horizontal Velocity



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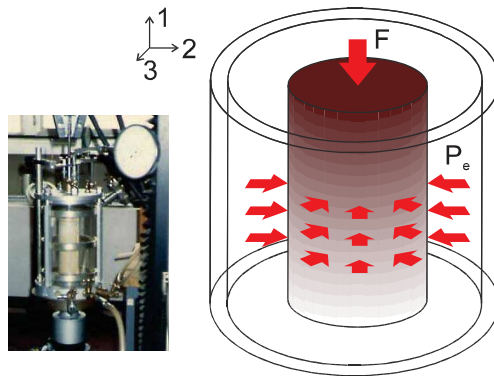
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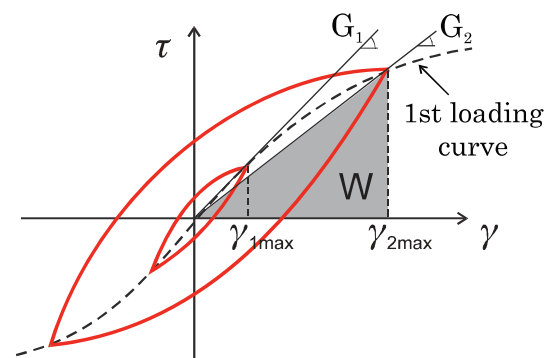
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■ Cyclic triaxial test



Hysteresis loops



■ Damping:

$$\beta = \frac{\Delta E}{4\pi W}$$

■ ΔE : energy dissipated in a cycle (area of the hysteresis loop)

■ W : elastic energy stored during 1/4 cycle

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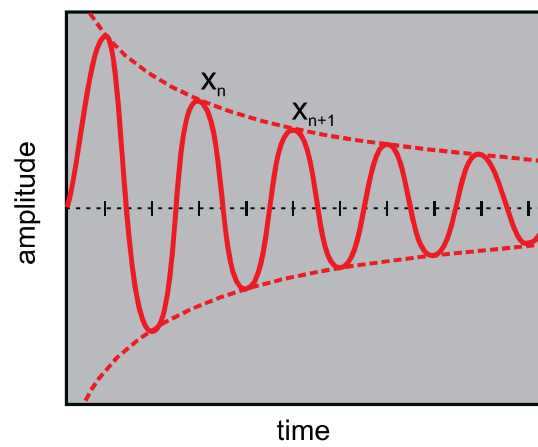
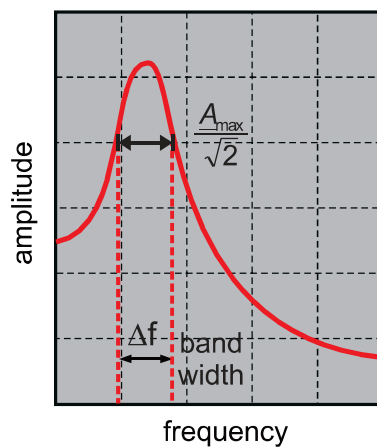
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■ Resonant column: direct estimation of damping:

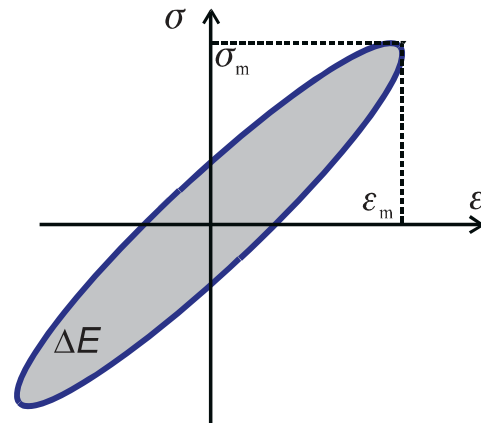


■ In freq. domain: $\beta = \frac{\Delta f}{2f_0}$

■ In time domain: $\delta = 2\pi\beta = \log \frac{x_{n+1}}{x_n}$

- Needs: damping for continuous models
- Framework: linear viscoelasticity
- Behaviour depends on **loading history**: $\underline{\underline{\sigma}} = f(\underline{\underline{\varepsilon}}, \underline{\underline{\dot{\varepsilon}}})$
 $\sigma = f(\varepsilon, \dot{\varepsilon})$ (longitudinal) or $\tau = f(\gamma, \dot{\gamma})$ (torsion), etc.

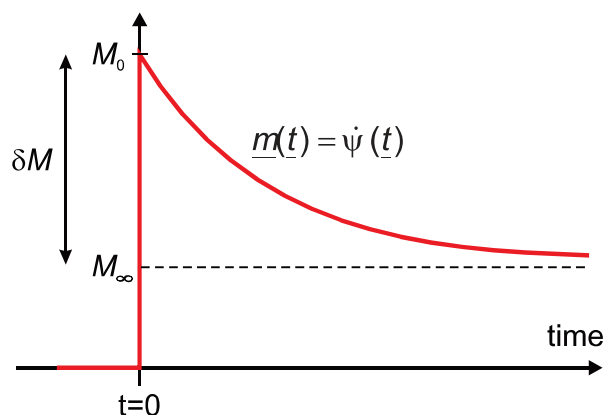
- Phase shift between stress and strain.
For $\hat{\varepsilon}(\omega) = \varepsilon_m \exp(i\omega t)$, the stress-strain curve is:



- The constitutive law may be written:

$$\sigma = \int_{-\infty}^{+\infty} m(t - \tau) \varepsilon(\tau) d\tau = m(t) * \varepsilon(t)$$

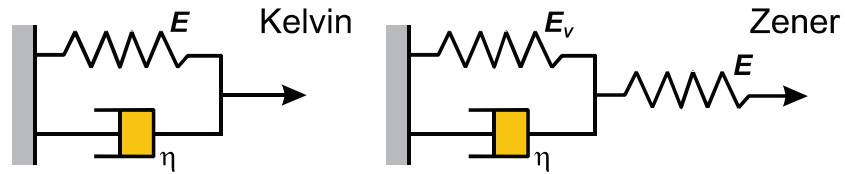
- $m(t)$: response in stress to a Dirac strain $\varepsilon(t) = \delta(t)$



M_0 : instantaneous modulus, M_∞ : relaxed modulus

- In the frequency domain, the constitutive law becomes:

$$\sigma(t) = m(t) * \varepsilon(t) \rightarrow \hat{\sigma}(\omega) = M(\omega)\hat{\varepsilon}(\omega)$$



- Kelvin-Voigt: $\sigma = \sigma_e + \sigma_v = E\varepsilon + \eta\dot{\varepsilon}$

$$\text{In freq.: } \hat{\sigma} = (E + i\omega\eta)\hat{\varepsilon} = M(\omega)\hat{\varepsilon}$$

- Zener: $\hat{\sigma} = \frac{E(E_v + i\omega\eta)}{E + E_v + i\omega\eta}\hat{\varepsilon} = \frac{1 + i\omega\tau_\varepsilon}{1 + i\omega\tau_\sigma}M_\infty\hat{\varepsilon}$

- Attenuation: $Q^{-1}(\omega) = \frac{M_I(\omega)}{M_R(\omega)} \simeq 2\beta(\omega) \quad Q_K^{-1}(\omega) = \frac{\omega\eta}{E}$

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- Equation of motion (freq.domain):

$$\frac{\partial \hat{\sigma}(x, \omega)}{\partial x} = \rho \frac{\partial^2 \hat{u}(x, \omega)}{\partial t^2} = -\rho \omega^2 \hat{u}(x, \omega)$$

- Viscoelastic behaviour: $\hat{\sigma}(x, \omega) = M(\omega) \hat{\varepsilon}(x, \omega)$

$$\frac{\partial \hat{\sigma}(x, \omega)}{\partial x} = M(\omega) \frac{\partial^2 \hat{u}(x, \omega)}{\partial x^2} = -\rho \omega^2 \hat{u}(x, \omega)$$

- Wave equation in the damped case:

$$\frac{\partial^2 \hat{u}(x, \omega)}{\partial x^2} = -\frac{\rho \omega^2}{M(\omega)} \hat{u}(x, \omega)$$

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- Solution: longitudinal damped wave

$$\hat{u}(x, \omega) = U_0(\omega) \exp[-i\hat{k}(\omega)x]$$

where $\hat{k}(\omega)$ is the complex wavenumber such as:

$$\hat{k}^2(\omega) = \frac{\rho \omega^2}{M(\omega)}$$

- Denoting $\hat{k}(\omega) = k(\omega) - i\alpha(\omega)$, the solution reads:

$$\hat{u}(x, \omega) = U_0(\omega) \exp[-ik(\omega)x] \exp[-\alpha(\omega)x]$$

propagation attenuation

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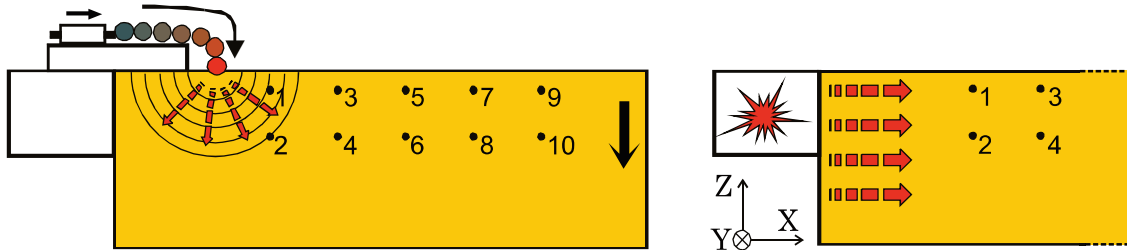
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Centrifuge experiments:



■ Solution 1: $\hat{u}(r, \omega) = \frac{A}{r} \left(-\frac{1}{r} - i\hat{k}(\omega) \right) \exp \left[-i\hat{k}(\omega)r \right]$

■ Solution 2: $\hat{u}(x, \omega) = A \exp \left[-i\hat{k}(\omega)x \right]$

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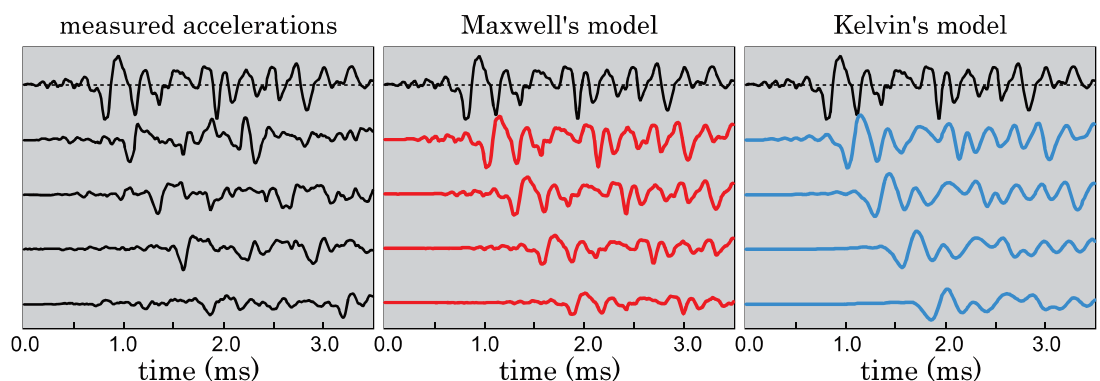
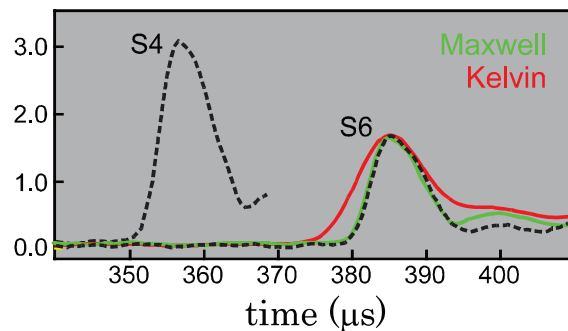
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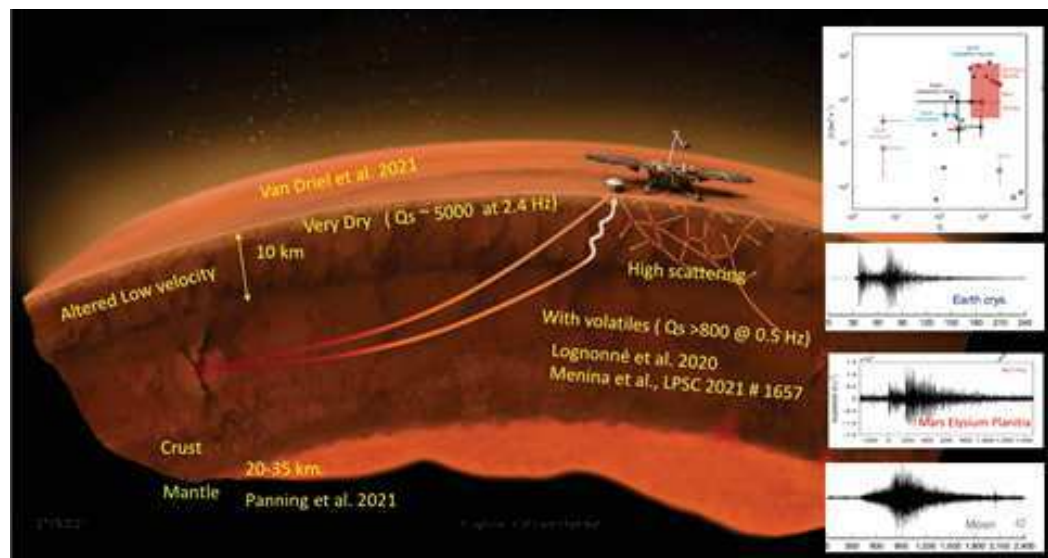
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- Observations from the Insight mission
- Mars crustal structure in Elysium Planitia
(Lognonné et al. 2020, Knapmeyer-Endrun et al., 2021)



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■ Strong motion: the Tohoku event



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